

ESTIMATES FOR THE VOLUME OF A LORENTZIAN MANIFOLD

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ABSTRACT. We prove new estimates for the volume of a Lorentzian manifold and show especially that cosmological spacetimes with crushing singularities have finite volume.

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0. INTRODUCTION

Let N be a $(n + 1)$ -dimensional Lorentzian manifold and suppose that N can be decomposed in the form

$$(0.1) \quad N = N_0 \cup N_- \cup N_+,$$

where N_0 has finite volume and N_- resp. N_+ represent the critical past resp. future Cauchy developments with not necessarily a priori bounded volume. We assume that N_+ is the future Cauchy development of a Cauchy hypersurface M_1 , and N_- the past Cauchy development of a hypersurface M_2 , or, more precisely, we assume the existence of a time function x^0 , such that

$$(0.2) \quad \begin{aligned} N_+ &= x^{0^{-1}}([t_1, T_+)), & M_1 &= \{x^0 = t_1\}, \\ N_- &= x^{0^{-1}}((T_-, t_2]), & M_2 &= \{x^0 = t_2\}, \end{aligned}$$

and that the Lorentz metric can be expressed as

$$(0.3) \quad d\bar{s}^2 = e^{2\psi} \{-dx^{0^2} + \sigma_{ij}(x^0, x) dx^i dx^j\},$$

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where $x = (x^i)$ are local coordinates for the space-like hypersurface M_1 if N_+ is considered resp. M_2 in case of N_- .

The coordinate system $(x^\alpha)_{0 \leq \alpha \leq n}$ is supposed to be future directed, i.e. the *past* directed unit normal (ν^α) of the level sets

$$(0.4) \quad M(t) = \{x^0 = t\}$$

is of the form

$$(0.5) \quad (\nu^\alpha) = -e^{-\psi}(1, 0, \dots, 0).$$

If we assume the mean curvature of the slices $M(t)$ with respect to the past directed normal—cf. [5, Section 2] for a more detailed explanation of our conventions—is strictly bounded away from zero, then, the following volume estimates can be proved

Theorem 0.1. *Suppose there exists a positive constant ϵ_0 such that*

$$(0.6) \quad H(t) \geq \epsilon_0 \quad \forall t_1 \leq t < T_+,$$

and

$$(0.7) \quad H(t) \leq -\epsilon_0 \quad \forall T_- < t \leq t_2,$$

then

$$(0.8) \quad |N_+| \leq \frac{1}{\epsilon_0} |M(t_1)|,$$

and

$$(0.9) \quad |N_-| \leq \frac{1}{\epsilon_0} |M(t_2)|.$$

These estimates also hold locally, i.e. if $E_i \subset M(t_i)$, $i = 1, 2$, are measurable subsets and E_1^+, E_2^- the corresponding future resp. past directed cylinders, then,

$$(0.10) \quad |E_1^+| \leq \frac{1}{\epsilon_0} |E_1|,$$

and

$$(0.11) \quad |E_2^-| \leq \frac{1}{\epsilon_0} |E_2|.$$

1. PROOF OF THEOREM 0.1

In the following we shall only prove the estimate for N_+ , since the other case N_- can easily be considered as a future development by reversing the time direction.

Let $x = x(\xi)$ be an embedding of a space-like hypersurface and (ν^α) be the past directed normal. Then, we have the Gauß formula

$$(1.1) \quad x_{ij}^\alpha = h_{ij} \nu^\alpha.$$

where (h_{ij}) is the second fundamental form, and the Weingarten equation

$$(1.2) \quad \nu_i^\alpha = h_i^k x_k^\alpha.$$

We emphasize that covariant derivatives, indicated simply by indices, are always *full* tensors.

The slices $M(t)$ can be viewed as special embeddings of the form

$$(1.3) \quad x(t) = (t, x^i),$$

where (x^i) are coordinates of the *initial* slice $M(t_1)$. Hence, the slices $M(t)$ can be considered as the solution of the evolution problem

$$(1.4) \quad \dot{x} = -e^\psi \nu, \quad t_1 \leq t < T_+,$$

with initial hypersurface $M(t_1)$, in view of (0.5).

From the equation (1.4) we can immediately derive evolution equations for the geometric quantities g_{ij}, h_{ij}, ν , and $H = g^{ij} h_{ij}$ of $M(t)$, cf. e.g. [3, Section 4], where the corresponding evolution equations are derived in Riemannian space.

For our purpose, we are only interested in the evolution equation for the metric, and we deduce

$$(1.5) \quad \dot{g}_{ij} = \langle \dot{x}_i, x_j \rangle + \langle x_i, \dot{x}_j \rangle = -2e^\psi h_{ij},$$

in view of the Weingarten equation.

Let $g = \det(g_{ij})$, then,

$$(1.6) \quad \dot{g} = g g^{ij} \dot{g}_{ij} = -2e^\psi H g,$$

and thus, the volume of $M(t)$, $|M(t)|$, evolves according to

$$(1.7) \quad \frac{d}{dt} |M(t)| = \int_{M(t_1)} \frac{d}{dt} \sqrt{g} = - \int_{M(t)} e^\psi H,$$

where we shall assume without loss of generality that $|M(t_1)|$ is finite, otherwise, we replace $M(t_1)$ by an arbitrary measurable subset of $M(t_1)$ with finite volume.

Now, let $T \in [t_1, T_+)$ be arbitrary and denote by $Q(t_1, T)$ the cylinder

$$(1.8) \quad Q(t_1, T) = \{ (x^0, x) : t_1 \leq x^0 \leq T \},$$

then,

$$(1.9) \quad |Q(t_1, T)| = \int_{t_1}^T \int_M e^\psi,$$

where we omit the volume elements, and where, $M = M(x^0)$.

By assumption, the mean curvature H of the slices is bounded from below by ϵ_0 , and we conclude further, with the help of (1.7),

$$(1.10) \quad \begin{aligned} |Q(t_1, T)| &\leq \frac{1}{\epsilon_0} \int_{t_1}^T \int_M e^\psi H \\ &= \frac{1}{\epsilon_0} \{ |M(t_1)| - |M(T)| \} \\ &\leq \frac{1}{\epsilon_0} |M(t_1)|. \end{aligned}$$

Letting T tend to T_+ gives the estimate for $|N_+|$.

To prove the estimate (0.10), we simply replace $M(t_1)$ by E_1 .

If we relax the conditions (0.6) and (0.7) to include the case $\epsilon_0 = 0$, a volume estimate is still possible.

Theorem 1.1. *If the assumptions of Theorem 0.1 are valid with $\epsilon_0 = 0$, and if in addition the length of any future directed curve starting from $M(t_1)$ is bounded by a constant γ_1 and the length of any past directed curve starting from $M(t_2)$ is bounded by a constant γ_2 , then,*

$$(1.11) \quad |N_+| \leq \gamma_1 |M(t_1)|$$

and

$$(1.12) \quad |N_-| \leq \gamma_2 |M(t_2)|.$$

Proof. As before, we only consider the estimate for N_+ .

From (1.6) we infer that the volume element of the slices $M(t)$ is decreasing in t , and hence,

$$(1.13) \quad \sqrt{g(t)} \leq \sqrt{g(t_1)} \quad \forall t_1 \leq t.$$

Furthermore, for fixed $x \in M(t_1)$ and $t > t_1$

$$(1.14) \quad \int_{t_1}^t e^\psi \leq \gamma_1$$

because the left-hand side is the length of the future directed curve

$$(1.15) \quad \gamma(\tau) = (\tau, x) \quad t_1 \leq \tau \leq t.$$

Let us now look at the cylinder $Q(t_1, T)$ as in (1.8) and (1.9). We have

$$(1.16) \quad \begin{aligned} |Q(t_1, T)| &= \int_{t_1}^T \int_{M(t_1)} e^\psi \sqrt{g(t, x)} \leq \int_{t_1}^T \int_{M(t_1)} e^\psi \sqrt{g(t_1, x)} \\ &\leq \gamma_1 \int_{M(t_1)} \sqrt{g(t_1, x)} = \gamma_1 |M(t_1)| \end{aligned}$$

by applying Fubini's theorem and the estimates (1.13) and (1.14). \square

2. COSMOLOGICAL SPACETIMES

A cosmological spacetime is a globally hyperbolic Lorentzian manifold N with compact Cauchy hypersurface \mathcal{S}_0 , that satisfies the timelike convergence condition, i.e.

$$(2.1) \quad \bar{R}_{\alpha\beta} \nu^\alpha \nu^\beta \geq 0 \quad \forall \langle \nu, \nu \rangle = -1.$$

If there exist crushing singularities, see [1] or [2] for a definition, then, we proved in [2] that N can be foliated by spacelike hypersurfaces $M(\tau)$ of constant mean curvature τ , $-\infty < \tau < \infty$,

$$(2.2) \quad N = \bigcup_{0 \neq \tau \in \mathbb{R}} M(\tau) \cup \mathcal{C}_0,$$

where \mathcal{C}_0 consists either of a single maximal slice or of a whole continuum of maximal slices in which case the metric is stationary in \mathcal{C}_0 . But in any case \mathcal{C}_0 is a compact subset of N .

In the complement of \mathcal{C}_0 the mean curvature function τ is a regular function with non-vanishing gradient that can be used as a new time function, cf. [4] for a simple proof.

Thus, the Lorentz metric can be expressed in Gaussian coordinates (x^α) with $x^0 = \tau$ as in (0.3). We choose arbitrary $\tau_2 < 0 < \tau_1$ and define

$$(2.3) \quad \begin{aligned} N_0 &= \{ (\tau, x) : \tau_2 \leq \tau \leq \tau_1 \}, \\ N_- &= \{ (\tau, x) : -\infty < \tau \leq \tau_2 \}, \\ N_+ &= \{ (\tau, x) : \tau_1 \leq \tau < \infty \}. \end{aligned}$$

Then, N_0 is compact, and the volumes of N_-, N_+ can be estimated by

$$(2.4) \quad |N_+| \leq \frac{1}{\tau_1} |M(\tau_1)|,$$

and

$$(2.5) \quad |N_-| \leq \frac{1}{|\tau_2|} |M(\tau_2)|.$$

Hence, we have proved

Theorem 2.1. *A cosmological spacetime N with crushing singularities has finite volume.*

Remark 2.2. Let N be a spacetime with compact Cauchy hypersurface and suppose that a subset $N_- \subset N$ is foliated by constant mean curvature slices $M(\tau)$ such that

$$(2.6) \quad N_- = \bigcup_{0 < \tau \leq \tau_2} M(\tau)$$

and suppose furthermore, that $x^0 = \tau$ is a time function—which will be the case if the timelike convergence condition is satisfied—so that the metric can be represented in Gaussian coordinates (x^α) with $x^0 = \tau$.

Consider the cylinder $Q(\tau, \tau_2) = \{\tau \leq x^0 \leq \tau_2\}$ for some fixed τ . Then,

$$(2.7) \quad |Q(\tau, \tau_2)| = \int_\tau^{\tau_2} \int_M e^\psi = \int_\tau^{\tau_2} H^{-1} \int_M H e^\psi,$$

and we obtain in view of (1.7)

$$(2.8) \quad \tau_2^{-1} \{|M(\tau)| - |M(\tau_2)|\} \leq |Q(\tau, \tau_2)|,$$

and conclude further

$$(2.9) \quad \lim_{\tau \rightarrow 0} |M(\tau)| \leq \tau_2 |N_-| + |M(\tau_2)|,$$

i.e.

$$(2.10) \quad \lim_{\tau \rightarrow 0} |M(\tau)| = \infty \implies |N_-| = \infty.$$

3. THE RIEMANNIAN CASE

Suppose that N is a Riemannian manifold that is decomposed as in (0.1) with metric

$$(3.1) \quad d\bar{s}^2 = e^{2\psi} \{ dx^{0^2} + \sigma_{ij}(x^0, x) dx^i dx^j \}.$$

The Gauß formula and the Weingarten equation for a hypersurface now have the form

$$(3.2) \quad x_{ij}^\alpha = -h_{ij} \nu^\alpha,$$

and

$$(3.3) \quad \nu_i^\alpha = h_i^k x_k^\alpha.$$

As default normal vector—if such a choice is possible—we choose the outward normal, which, in case of the coordinate slices $M(t) = \{x^0 = t\}$ is given by

$$(3.4) \quad (\nu^\alpha) = e^{-\psi}(1, 0, \dots, 0).$$

Thus, the coordinate slices are solutions of the evolution problem

$$(3.5) \quad \dot{x} = e^\psi \nu,$$

and, therefore,

$$(3.6) \quad \dot{g}_{ij} = 2e^\psi h_{ij},$$

i.e. we have the opposite sign compared to the Lorentzian case leading to

$$(3.7) \quad \frac{d}{dt} |M(t)| = \int_M e^\psi H.$$

The arguments in Section 1 now yield

Theorem 3.1. (i) *Suppose there exists a positive constant ϵ_0 such that the mean curvature $H(t)$ of the slices $M(t)$ is estimated by*

$$(3.8) \quad H(t) \geq \epsilon_0 \quad \forall t_1 \leq t < T_+,$$

and

$$(3.9) \quad H(t) \leq -\epsilon_0 \quad \forall T_- < t \leq t_2,$$

then

$$(3.10) \quad |N_+| \leq \frac{1}{\epsilon_0} \lim_{t \rightarrow T_+} |M(t)|,$$

and

$$(3.11) \quad |N_-| \leq \frac{1}{\epsilon_0} \lim_{t \rightarrow T_-} |M(t)|.$$

(ii) *On the other hand, if the mean curvature H is negative in N_+ and positive in N_- , then, we obtain the same estimates as Theorem 0.1, namely,*

$$(3.12) \quad |N_+| \leq \frac{1}{\epsilon_0} |M(t_1)|,$$

and

$$(3.13) \quad |N_-| \leq \frac{1}{\epsilon_0} |M(t_2)|.$$

REFERENCES

- [1] D. Eardley & L. Smarr, *Time functions in numerical relativity: marginally bound dust collapse*, Phys. Rev. D **19** (1979) 2239–2259. 5
- [2] C. Gerhardt, *H-surfaces in Lorentzian manifolds*, Commun. Math. Phys. **89** (1983) 523–553. 5
- [3] ———, *Hypersurfaces of prescribed Weingarten curvature*, Math. Z. **224** (1997) 167–194. gerhardt/preprints 3
- [4] ———, *On the foliation of space-time by constant mean curvature hypersurfaces*, e-print, 7 pages, math.DG/0304423 5
- [5] ———, *Hypersurfaces of prescribed curvature in Lorentzian manifolds*, Indiana Univ. Math. J. **49** (2000) 1125–1153. gerhardt/preprints 2
- [6] S. W. Hawking & G. F. R. Ellis, *The large scale structure of space-time*, Cambridge University Press, Cambridge, 1973.

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